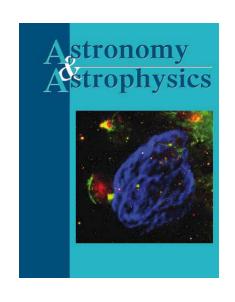


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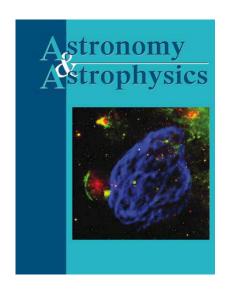
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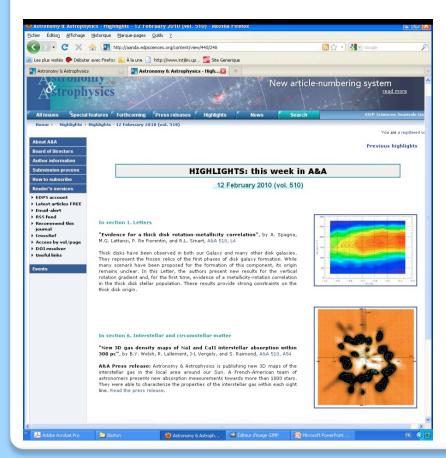


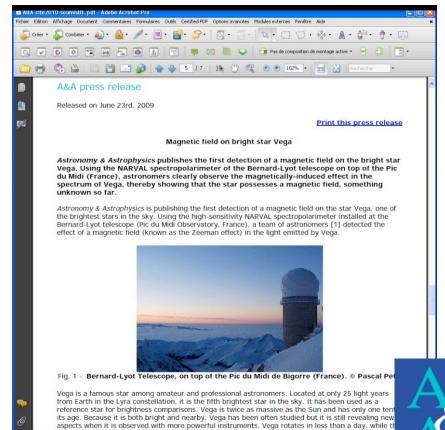
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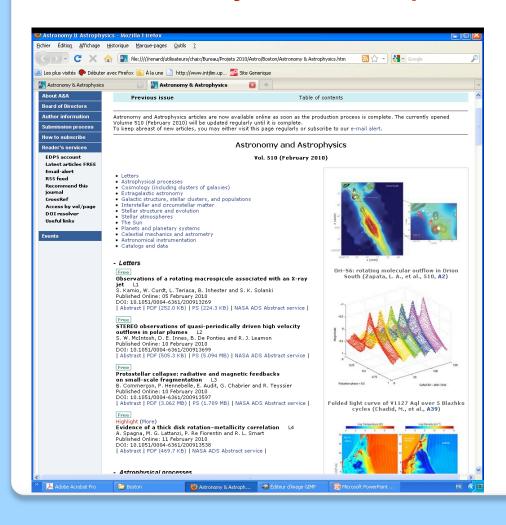
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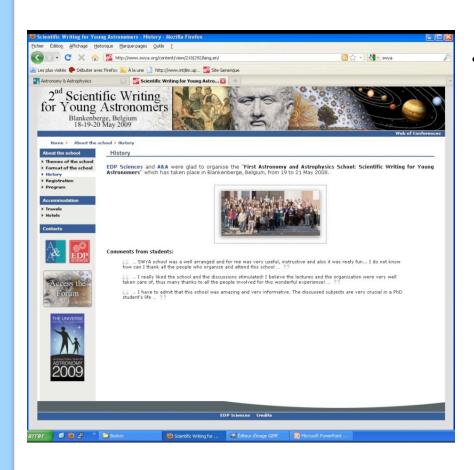


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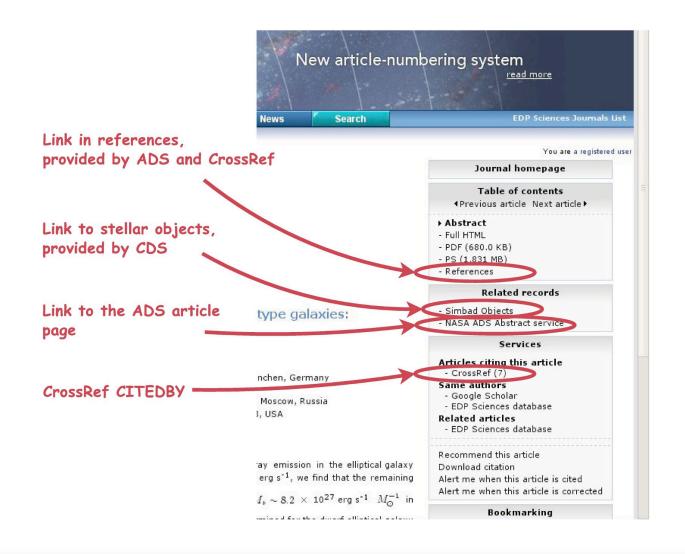


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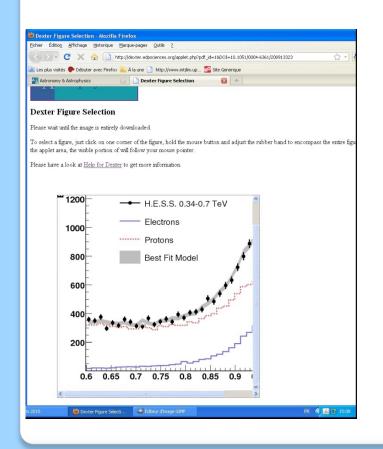










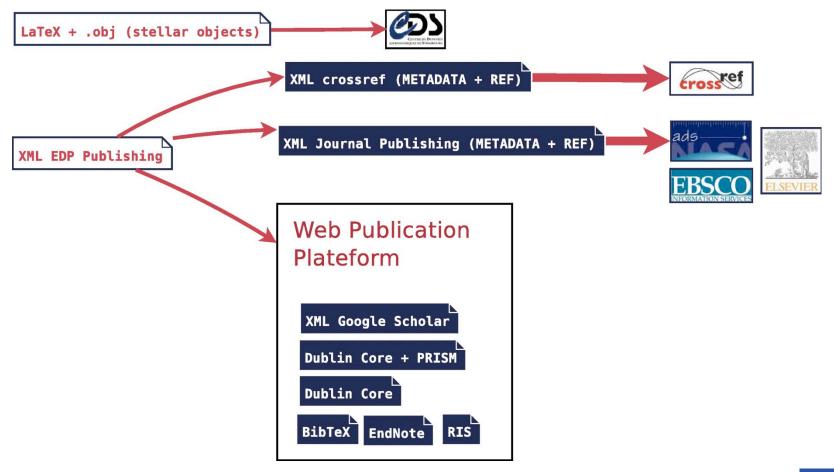


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Example:

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Citation

Jutzi, M., Michel, P., & Benz, W. 2010, A&A, 509, L2 Janiuk, A., & Yuan, Y-F., 2010, A&A, 509, A55



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Astronomy and Astrophysics: electronic first

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Astronomy Astrophysics

The role of black hole spin and magnetic field threading the unstable neutrino disk in gamma ray bursts

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tori have been already discussed in a number of works (e.g. Popham et al. 1999; Di Matteo et al. 2002; Kohri et al. 2005; Chen & Beloborodov 2007; Janiuk et al. 2004, 2007; Lei et al.

we develop here was first described in Janiuk et al. (2004) for rale and viscosity in the disk. We also check what effects may be

plausible mechanism that helps to launch the jets emitted in a narrow cone along the rotation axis. As shown in Janiuk et al. (2008), especially for the long duration GRBs the BH spin is required to sustain the central engine activity and the jet produc-In the present paper we expand on our previous work, and we tion for a sufficiently long time. Because these long GRBs offurther investigate the properties and evolution of such hot and ten exhibit a very variable temporal structure (e.g. Beloborodov dense accretion tori. At the extreme densities and temperatures, et al. 2000), it seems very important to study a possible instadetermined by the hyper-Eddington accretion rate, the torus is billity mechanism that may account for the variability. We check cooled mainly by neutrino emission produced primarily by elec-here for which values of the black hole spin the instability may tron and positron capture on nucleons (6 reactions). The model operate near the central black hole, depending on the accretion

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imposed on the torus and its unstable strip by the additional heat- and from protons to neutrons for a given buryon number dening via the magnetic fields, threaded by the rotating black hole. The content of the article is as follows. First, we discuss the busic assumptions and present the main equations of the model. that are the result of a black hole rotation with an arbitrary spin. charge neutrality (Yuan 2005). Third, we introduce the energy extraction from the rotating black hole via the magnetic field, as an additional physical process that present the results of our calculations, and in Sect. 4 we discuss the results and conclude.

2 Model

The model of a neutrino-cooled accreting torus was fully discussed in Janiuk et al. (2004, 2007). Here we briefly repeat the basic assumptions, and the model equations are given in Sect. 2.1. In Sect. 2.2 we describe the relativistic corrections to

sity, n_b , and temperature, T. These reactions are electron and positron capture on nucleons and neutron decay (see Kohri et al. 2005; Janiuk et al. 2007). The closing equations for the EOS are Second, we introduce the changes and corrections to the model the conservation of the baryon number, $n_0 + n_0 = n_0 X_{\text{nuc}}$, and the For the neutrino cooling of the torus, we consider the

electron-positron pair annihilation, bremsstrahlung, plasmon demay operate in the gamma ray burst central engine. In Sect. 3 we cay and beta reactions. Each of these emission processes has a reverse one, which leads to neutrino absorption, and we calculate the absorptive optical depth ran, for the neutrinos of the three flavors. In addition, the free escape of neutrinos from the disc is limited by scattering, and we calculate the scattering optical depth $\tau_{\rm D}$. The neutrino cooling rate is therefore in the neutrinothick torus given by

$$Q_r^* = \frac{\frac{1}{2}\sigma T^4}{\frac{1}{2}} \sum_{\underline{\alpha_0 + \epsilon_1} + \underline{1}_{\underline{\alpha}} + \underline{1}_{\underline{\alpha}}} \frac{1}{1}$$
 (3)

the fautation pressure multures are factor use to are electron positron pairs, and the neutrino pressure is calculated for the trapped neutrinos, using the two-stream approximation (Popham & Narayan 1995; Di Matteo et al. 2002). We consider both the neutrino transparent and opaque regions as well as the transition between the two, and in order to determine the distribution function of the partially trapped neutrinos we use a "gray body" model (Sawyer 2003), consistent with the two-stream approximation.

The chemical potentials, or equivalently the ratio of free protons, $x = n_p/n_b$, are determined from the condition of equilibrium between the transition reactions from neutrons to protons $F_{tot} = Q_{visc}^* = Q_{adv}^* + Q_{cad}^* + Q_{c}^* + Q_{photo-}$

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The total flux of energy generated at the radius r is determined by the global model parameters, i.e. the black hole mass and accretion rate:

$$F_{tot} = \frac{3GMM}{8\pi r^3} f(r) \qquad (9)$$

where f(r) stands for the inner boundary condition and will be given below. In order to calculate the initial stationary configuration, we solve the energy balance:

$$F_{tot} = Q_{visc}^{+} = Q_{adv}^{-} + Q_{rad}^{-} + Q_{v}^{-} + Q_{photo}$$
. (10)





Consequences

- Articles are online earlier, around 10-15 days
- On the tools:

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Bibcode: 2010A&A...510A...2Z

Bibtex:

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volume = 510,

pages = {A2-+},

...
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